

The energy density in the  
2D Ising model.

Clément Hongler.  
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Joint work with Stas Smirnov.

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The Ising model on a  
graph  $G$

• Probability space:  $\{\pm 1\}^{V(G)}$

• For  $\vec{\sigma} \in \{\pm 1\}^{V(G)}$ ,

$$\mathbb{P}[\vec{\sigma}] = \frac{1}{Z_\beta} e^{-\beta H(\vec{\sigma})}$$

•  $\beta > 0$  (inverse temperature)

•  $H(\vec{\sigma}) = -\sum_{\langle x,y \rangle \in E(G)} \sigma_x \sigma_y$  (energy)

$$Z_\beta = \sum_{\vec{\sigma} \in \{\pm 1\}^{V(G)}} e^{-\beta H(\vec{\sigma})}$$

(the partition function)

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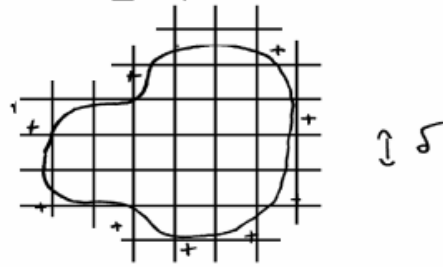
$$H(\vec{\sigma}) = - \sum_{\langle x, y \rangle \in E(G)} \sigma_x \sigma_y \quad (\text{energy})$$

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The Ising model on the square grid.

•  $\Omega$  a Jordan domain.



$$\Omega^\delta := \Omega \cap \delta \mathbb{Z}^2$$

(largest connected component)

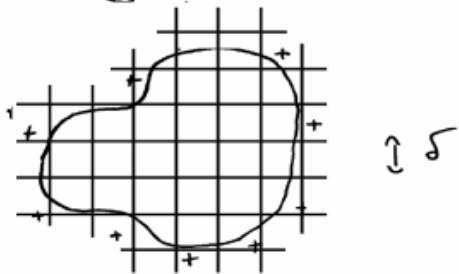
Ising model on  $\Omega^\delta$  with  
+ boundary condition

• What happens as  $\delta \rightarrow 0$ ?

The Ising model on the square grid.

It depends on  $\beta$ !  
 ( $\mathbb{P}[\vec{\sigma}] \propto e^{-\beta H(\vec{\sigma})}$ )

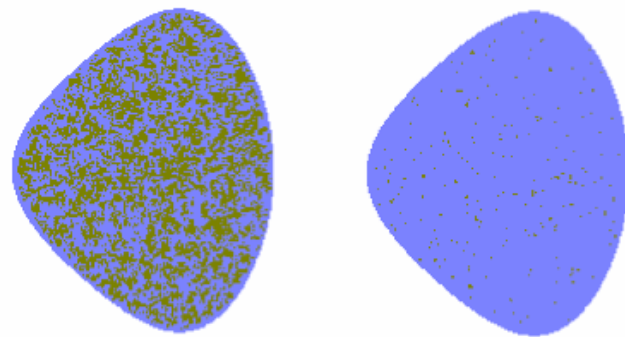
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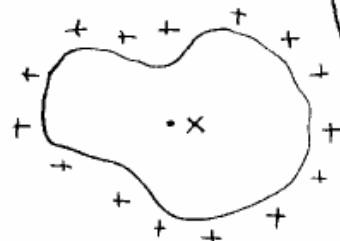
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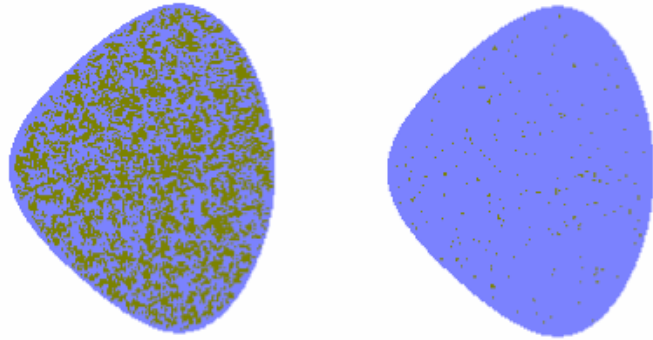
$$\beta < \beta_c = \frac{1}{2} \ln(\sqrt{2} + 1)$$

Thm (Onsager, 1944)

$$\mathbb{E}[\sigma_x] \xrightarrow{\delta \rightarrow 0} \begin{cases} 0 & \text{if } \beta \leq \beta_c \\ c(\beta) > 0 & \text{if } \beta > \beta_c \end{cases}$$



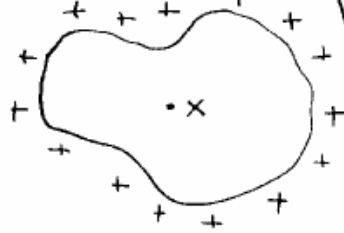
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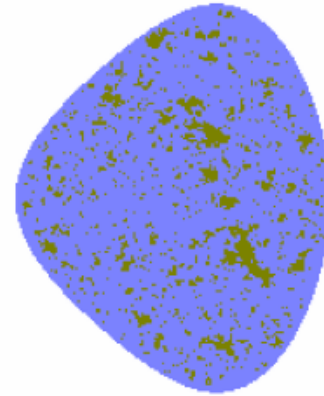
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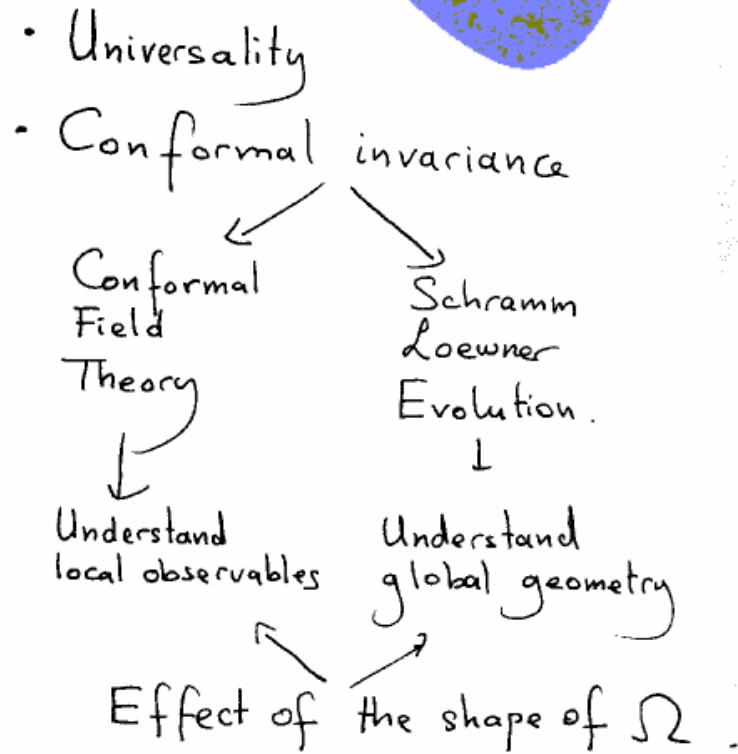
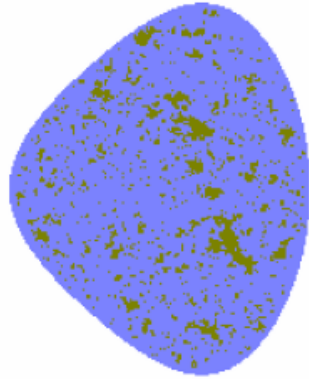


What happens at  $\beta_c$ ?



- Universality
  - Conformal invariance
    - Conformal Field Theory
      - Understand local observables
    - Schramm Loewner Evolution
      - Understand global geometry
- Effect of the shape of  $\Omega$ .

What happens at  $\beta_c$ ?



The energy density.

The energy terms in

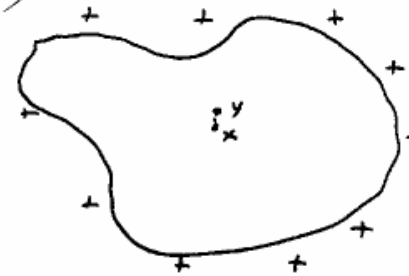
$$H(\vec{\sigma}) = - \sum_{\langle x, y \rangle \in E(G)} \sigma_x \sigma_y$$

Question: how does the energy distribute across the lattice? Fix an edge  $a = \langle x, y \rangle$ .

Thm (H., Smirnov)

$$\frac{P[\sigma_x = \sigma_y] - \frac{2+\sqrt{2}}{4}}{\delta}$$

$$\xrightarrow{\delta \rightarrow 0} \frac{1}{4\pi} \mathcal{L}_{\Omega}(a)$$



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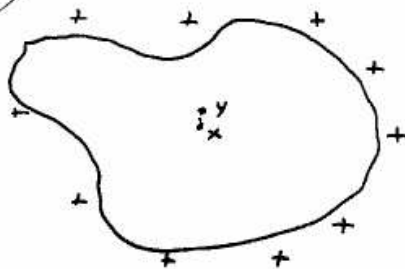
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The hyperbolic metric element

$$\mathcal{L}_{\mathcal{D}(a_1)}(a) = \frac{1}{1-|a|^2}$$



(M.C. Escher)

The energy density.

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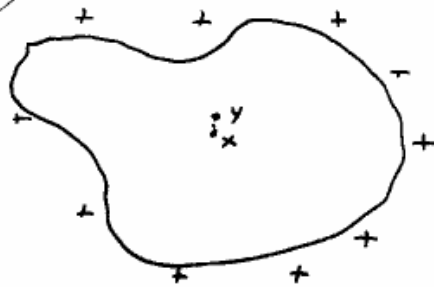
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Sketch of the proof.

(A) Contour representation

(B) Holomorphic deformation.

(C) Limit  $\delta \rightarrow 0$ .

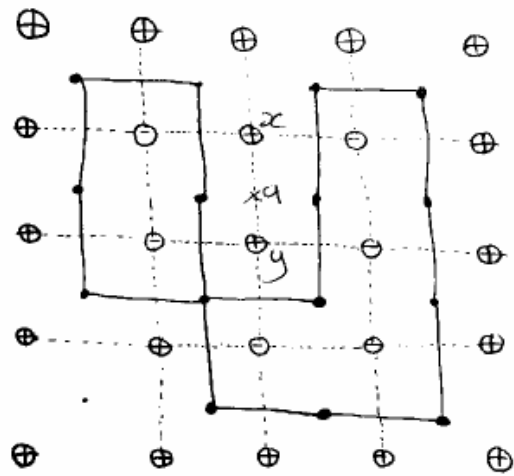


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Bijection  $\{\pm 1\}^{V(G)} \leftrightarrow \mathcal{C}$

$\mathcal{C} := \{\gamma \in E(G^*) : \gamma \text{ coll of closed loops}\}$

$$\mathbb{P}[\gamma] = \frac{1}{Z} \alpha^{\#\text{edges}(\gamma)}$$

$$\alpha = \alpha_c = \sqrt{2} - 1.$$

$$Z = \sum_{\gamma \in \mathcal{C}} \alpha^{\#\text{edges}(\gamma)}$$

Identify  $a$  to its midpoint

$$\mathcal{C}_a = \{\gamma \in \mathcal{C} : a \notin \gamma\}$$

$$Z_a = \sum_{\gamma \in \mathcal{C}_a} \alpha^{\#\text{edges}(\gamma)}$$

$$\mathbb{P}[\sigma_x = \sigma_y] = \frac{Z_a}{Z}.$$

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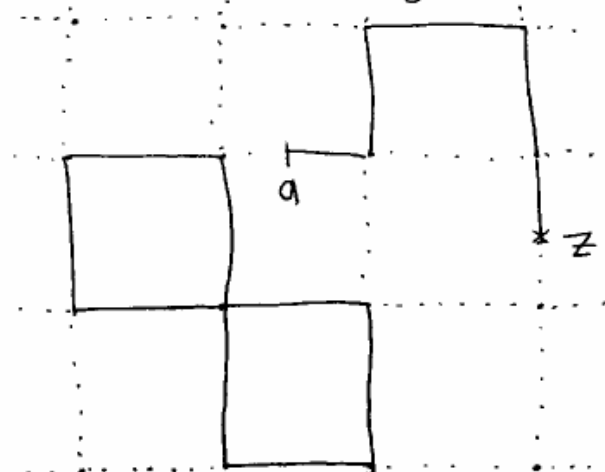
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Idea: construct a (discrete) holomorphic deformation of  $\frac{Z_a}{Z}$

$\mathcal{C}_{a,z} = \{\gamma \in E(G) : \gamma \text{ contains closed loops and a path from } a \text{ to } z \text{ starting from left to right at } a\}$

$z$ : midpoint of edge



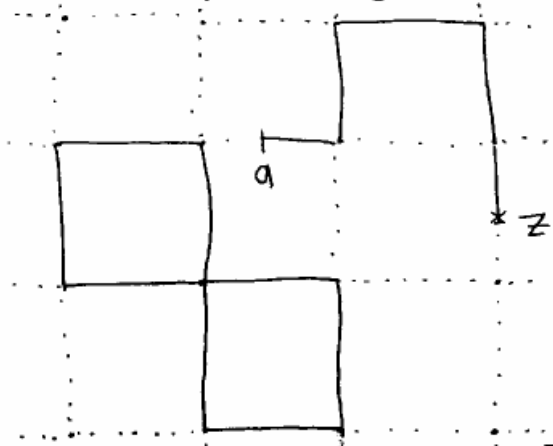
$$W(\gamma) = \frac{-\pi}{2}$$

# edges = 12

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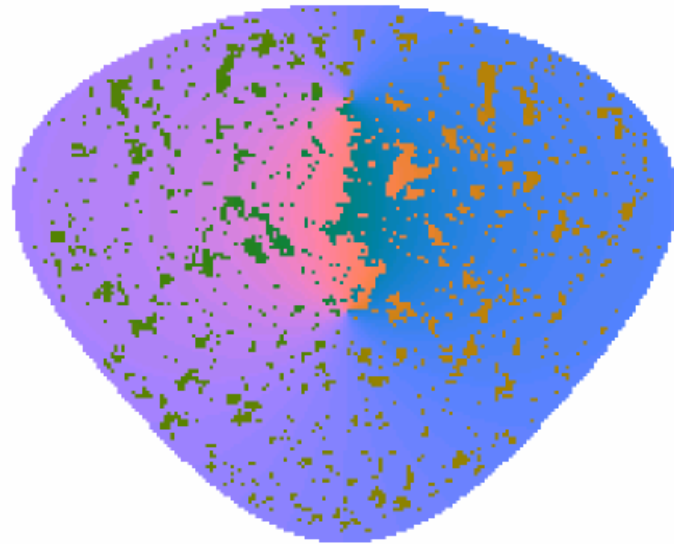
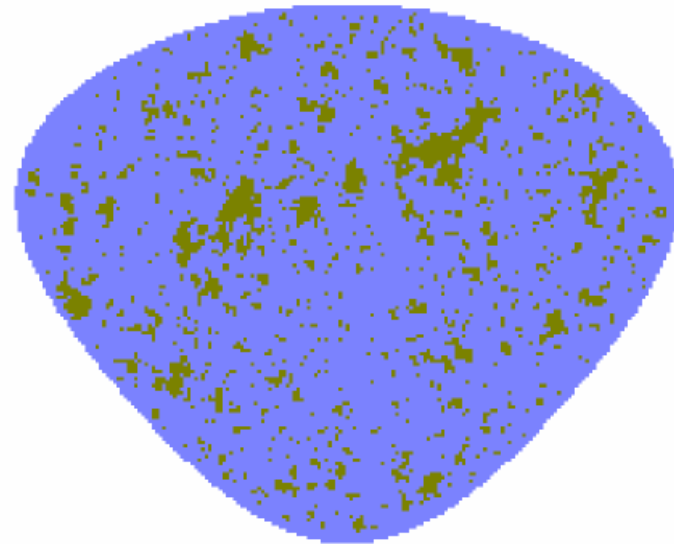
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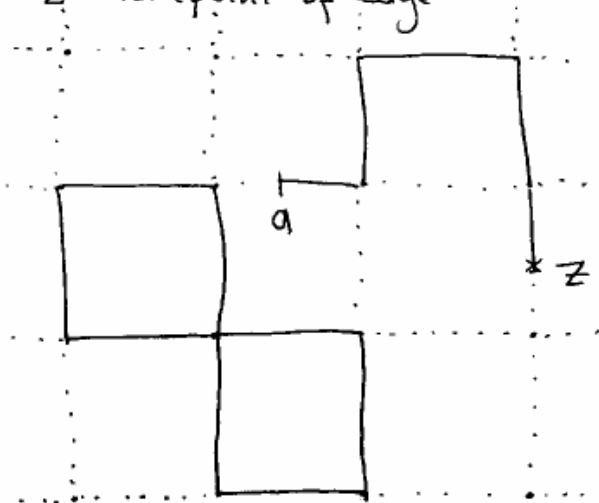
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$W$ : winding number of the path from  $a$  to  $z$ .

(total rotation of the path).

Holomorphic observable.

$$f_a^\sigma(z) := \frac{\sum_{\gamma \in C_{a,z}^\sigma} a^{\# \text{edges}(\gamma)} e^{-i \frac{W(\gamma)}{z}}}{z^\sigma}$$

$$f_a^\sigma(a) := \frac{z_a^\sigma}{z^\sigma}$$

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$$f_a^\delta(z) := \frac{\sum_{\gamma \in C_{a,z}^\delta} \alpha^{\#\text{edges}(\gamma)} e^{-i \frac{W(\gamma)}{2}}}{Z^\delta}$$

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Properties of  $z \mapsto f_a^\delta(z)$

(1) Discrete holomorphic on  $\Omega \setminus \{a\}$   
 $\Rightarrow \bar{\partial}^\delta f_a^\delta = 0, \bar{\partial} = \frac{1}{2}(\partial_x + i\partial_y)$

(2) Discrete singularity at  $a$ .  
 $(\bar{\partial}^\delta f_a^\delta)(a) = \frac{1}{2}$

(3) Boundary condition:  
 $f_a^\delta|_{\partial\Omega} \ll \frac{1}{\sqrt{n}}$ ;  $n$ : normal to ext.

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Derivation of the energy density

(1) + (2)  $\Leftrightarrow h_a^\delta := f_a^\delta - g_a^\delta$   
discrete analytic on  $\Omega_j$

$g_a^\delta$  discrete Green function

$$g_a^\delta(a) = \frac{\sqrt{2} + 2}{4}$$

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## Proof of convergence

1) Precompactness.  
(difficult: control  
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2) Identification of  
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 $f_a^\delta \parallel \frac{1}{\sqrt{h}}$   $(f_a^\delta)^2 \parallel \frac{1}{h}$   
 $\text{Re} \left( \int_{\Omega} (f_a^\delta)^c \right) \Big|_{\partial\Omega} = \text{cst}$   
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## Proof of precompactness

Assume  $\partial\Omega$  of finite length

Idea: control  $H_a^\delta := \operatorname{Re} \left( \int (h_a^\delta)^c \right)$

•  $H_a^\delta$  subharmonic

$$\partial_n H_a^\delta \leq M \text{ on } \partial\Omega.$$

$$0 \leq \int \int \Delta H_a^\delta = \int (\partial_n H_a^\delta)_+ - (\partial_n H_a^\delta)_-$$

$$\int (\partial_n H_a^\delta)_- \leq \int (\partial_n H_a^\delta)_+ \\ \leq M \ell(\partial\Omega)$$

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$$H_a^\delta = H + S, \text{ with}$$
$$\Delta H = 0, \quad \int \int_{\partial\Omega} = 0.$$
$$H|_{\partial\Omega} = H_a^\delta|_{\partial\Omega} \quad \Delta S \geq 0.$$

Control of  $S$  follows  
from  $\int \int \Delta S \leq M \ell(\partial\Omega)$

Control of  $H$  follows  
from  $\sum |\partial_n H| \leq 2M \ell(\partial\Omega)$   
Idea: pass to harmonic  
conjugate  $\leadsto \sum |\partial_t| \ll \ell \text{ bd}$   
 $\Rightarrow C \text{ bd.}$



$$H_a^\sigma = H + S, \text{ with}$$

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Thank you!